

Effect of Interactions on Vortices in a Nonequilibrium Polariton Condensate

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We demonstrate the creation of vortices in a macroscopically occupied polariton state formed in a semiconductor microcavity. A weak external laser beam carrying orbital angular momentum (OAM) is used to imprint a vortex on the condensate arising from the polariton optical parametric oscillator (OPO). The vortex core radius is found to decrease with increasing pump power, and is determined by polariton-polariton interactions. As a result of OAM conservation in the parametric scattering process, the excitation consists of a vortex in the signal and a corresponding antivortex in the idler of the OPO. The experimental results are in good agreement with a theoretical model of a vortex in the polariton OPO.

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Microcavity polaritons are bosonic quasiparticles, which arise from strong exciton-photon coupling in semiconductor microcavities. Characteristic bosonic phenomena, such as stimulated scattering [1] and polariton condensation have been reported [2]. In contrast to well-known atom Bose-Einstein condensates (BEC), polariton condensates are intrinsically out of equilibrium: a balance is reached between external excitation and losses due to the finite polariton lifetime [3,4]. In this respect the system has similarities to a laser, but with the fundamental difference that the particles have hybrid light-matter character with strong interactions due to the excitonic component.

Nonequilibrium dynamics can lead to a much richer phenomenology of spontaneous pattern formation than in equilibrium systems [5], with vortices a typical example. Vortices are found in many areas of science, including particle physics, optics, and condensed matter systems. They are characterized by winding of phase around a point, the vortex core, where the density of the system goes to zero. The total phase change for a complete loop must be an integer multiple of 2π , so a vortex is a state with quantized orbital angular momentum (OAM). Vortices have been intensively studied in classical systems including optical vortex beams [6] and lasers [7]. They are also a characteristic property of condensed-phase systems like BEC, liquid helium, and superconductors [8].

Spontaneous formation of vortices has been observed in an incoherently pumped polariton system [9]. In this Letter we demonstrate that vortices in a polariton condensate can be created using a weak external imprinting beam (which we term the *im* beam). We find that the vortex core radius is determined by polariton-polariton interactions and show that these lead to a decrease of the radius with increasing particle density. Here, the low polariton mass ($\sim 10^{-5}m_e$) plays an important role; the polariton vortex radius is typically $\sim 10 \mu\text{m}$, some 2 orders of magnitude larger than in an atom system. This enables us to observe *in*

situ the vortex profile and its density dependence, unlike in atom BEC experiments where vortices can only be imaged following condensate expansion [10], which is a destructive technique.

The system we study is a microcavity optical parametric oscillator (OPO) [1], in which coherent high density signal (*s*) and idler (*i*) states (condensates) arise from polariton-polariton scattering from the pump (*p*), as shown in Fig. 1(a). An important property of the OPO is that the phase of the signal and idler are not determined by that of the pump [11], but appear as a spontaneous symmetry breaking at the OPO threshold, similar to the phase transition of a thermal equilibrium BEC. Both the OPO modes and the polariton condensate which is realized by non-resonant pumping [2] exhibit characteristic properties of a system possessing a macroscopic wave function, such as extended spatial and temporal coherence [2,12–14].

As well as vortex creation, we demonstrate OAM conservation in the parametric scattering process; this means that when a vortex is imprinted onto the OPO signal, polariton-polariton scattering results in an antivortex with opposite OAM in the idler, as has been predicted previously theoretically [15]. It is this vortex-antivortex pair that is created in our experiment. We show theoretically that when the signal population is small compared to the pump, which is the regime of the present experiments, the signal and idler vortices have identical profiles, determined by the interactions.

The experiments were carried out at 4 K on a GaAs-based microcavity grown by molecular beam epitaxy. It has a Rabi splitting of 6 meV and zero detuning between the quantum well exciton and the optical mode. A single mode Ti-Al₂O₃ laser on resonance with the lower polariton branch at 14° was used to excite the OPO [Fig. 1(a)]. The zero OAM pump was focused to a $\sim 50 \mu\text{m}$ Gaussian spot. A spatial light modulator was used to generate a Gauss-Laguerre optical beam, corresponding to an optical vortex

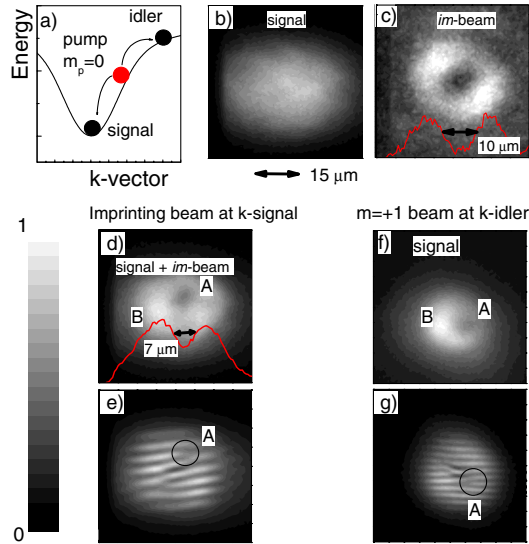


FIG. 1 (color online). (a) Schematic diagram of OPO. (b) Real space image of the signal with no imprinting beam. (c) Image of Gauss-Laguerre imprinting beam (*im* beam). (d) Signal with weak *im* beam of (c), showing an imprinted vortex labeled A. (e) Interferogram revealing the 2π phase variation around the vortex of (d). (f) Vortex in signal, labeled A, created by excitation at the idler position with Gauss-Laguerre beam. (g) Interferogram revealing the vortex with OAM $m_s = -1$ created in (f). The red lines are cross sections through the vortex cores of (c) to (d), with the sizes of the cores (FWHM) indicated.

with OAM $M = 1$ which was used as the imprinting beam.

Figure 1(b) shows the spatial image of the OPO signal emission, which occurs at 1.533 eV and has a linewidth of $<60 \mu\text{eV}$ (resolution limited), at a pump power of 3 times the OPO threshold, $P_{\text{th}} \sim 600 \text{ W cm}^{-2}$. It has a uniform spatial distribution, consistent with a sample having very weak disorder [16]. In order to create a vortex in the signal we introduce a weak *im* beam of power $\approx 10 \mu\text{W}$, ~ 40 times less than that of the signal, having OAM $M = +1$ and optical vortex core diameter of $10 \mu\text{m}$. The *im* beam is in resonance with the signal at normal incidence. Application of the *im* beam of Fig. 1(c) results in a strong modification of the spatial distribution of the signal [Fig. 1(d)], characterized by a resultant well-defined dip of diameter (FWHM) $\approx 7 \mu\text{m}$ [radius $\approx 3.5 \mu\text{m}$ (HWHM)] labeled A on Fig. 1(d).

To demonstrate that the dip arises from a vortex, the spatial phase variation of the signal was measured [9]. Figure 1(e) shows the interference pattern between the signal image in Fig. 1(d) and the image inverted around a central point of symmetry so that region A interferes with region B [see Fig. 1(d)], where the phase is nearly constant. The two forklike dislocations in the interference pattern demonstrate the presence of a single vortex, with phase winding by 2π about the core. One fork occurs at A, the position of the vortex in the original beam, while the other occurs at the equivalent location in the inverted beam.

We describe the process using the term “imprinting” because the *im* beam is very weak; we are imaging not the polaritons injected directly by the *im* beam, but the OPO signal modified by their presence. Indeed, the vortex disappears when the *im* beam is tuned off resonance with the signal. The reason imprinting can occur is that the phase of the signal is undetermined in the OPO, so even the few polaritons the *im* beam injects are able to lock the signal phase to their own. Experimentally there is a minimum ratio of the *im* beam to signal power density, of about $\sim 1/45$, required to create a vortex in the signal. This limit is probably determined by the fluctuations which lead to decoherence; [13,14] at lower powers the *im* beam is unable to overcome their effects.

The fact that the dip in the vortex core does not go fully to zero (the intensity typically reduces by a factor of ~ 1.6) suggests that the observed signal is not a simple OAM state with $m_s = 1$, but contains also some background with $m_s = 0$. This probably arises from unresolved multimode spectral structure in the signal, with the vortex only imprinted on the stronger features.

Vortices and OAM conservation have previously been investigated in nonlinear optics using down-conversion experiments [17]. However, the nonlinearity in those works has a $\chi^{(2)}$ form. A fundamental difference in the polariton system is that the nonlinearity is $\chi^{(3)}$. This means that, as well as giving rise to the OPO process, it leads to self- (and mutual) interaction terms, which give rise to the blueshifts of the modes [18]. These interaction terms are not found in $\chi^{(2)}$ processes. The polariton vortex profiles are expected to be determined by these interactions, with a characteristic length scale, the healing length, which depends on the population, as we show below.

In order to investigate this population dependence we have measured the vortex radius at different pump powers. Since the signal population increases with the pump power the *im*-beam power was adjusted to keep the ratio of the signal to *im*-beam power density constant at ~ 15 . Figures 2(a) and 2(b) show the vortex image at $P = 0.9 \text{ kW cm}^{-2} \sim 1.5P_{\text{th}}$ and $P = 2.6 \text{ kW cm}^{-2} \sim 4.5P_{\text{th}}$, respectively. The reduction of the vortex size with increasing excitation power is apparent. Figure 2(c) shows the dependence of the vortex core radius on excitation power; a decrease from $\sim 5.5 \mu\text{m}$ at threshold down to $\sim 3 \mu\text{m}$ at excitation density 5 times above threshold is observed. Moreover, very similar vortex sizes and variation with polariton density were obtained for two different sizes of the *im* beam with vortex core radius of $\sim 4 \mu\text{m}$ (circles) and $\sim 7 \mu\text{m}$ (triangles). These results show that the profile of the imprinted vortex is an intrinsic property of the interacting polariton system.

The theory below shows that the vortex radius is given by the healing length $\xi = (2M_c \kappa \sqrt{n_s n_i})^{-1/2}$, where M_c is the cavity photon effective mass, κ the strength of the nonlinearity, and n_s , n_i the signal and idler population densities. This can be understood qualitatively using the

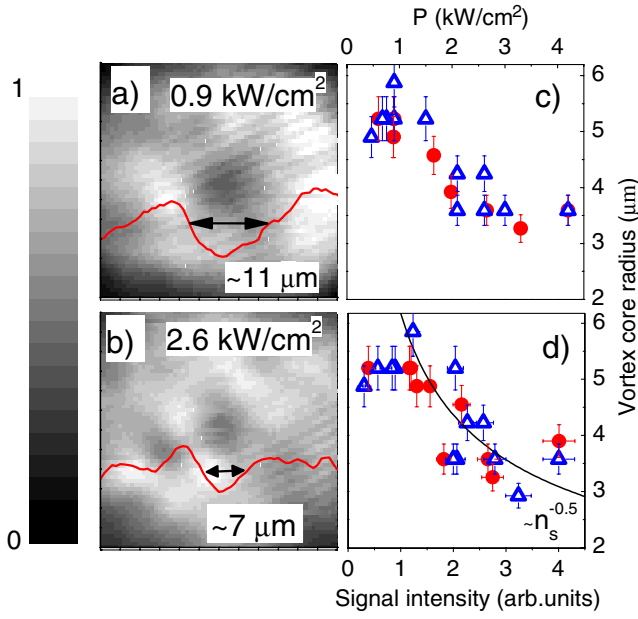


FIG. 2 (color online). (a),(b) Images of a signal vortex at (a) 0.9 kW cm^{-2} and (b) 2.6 kW cm^{-2} . Experimental vortex core radius of the OPO signal as a function of pump power (c) and as a function of the signal intensity (d) for an *im*-beam core radius of $\sim 4 \text{ } \mu\text{m}$ (circles) and $\sim 7 \text{ } \mu\text{m}$ (triangles), respectively.

following argument, which applies to interacting equilibrium BECs [8], and to lasers (VCSELs) with $\chi^{(3)}$ nonlinearity [7]: the healing length is obtained by equating a typical kinetic energy associated with a vortex in the condensate, $\sim (2M_c \xi^2)^{-1}$ to the interaction energy (blueshift), which is $\approx \kappa \sqrt{n_s n_i}$ (all energies are scaled to \hbar^2). Since n_s is proportional to n_i we expect the vortex radius to be approximately proportional to $n_s^{-1/2}$. Figure 2(d) shows that at higher signal intensities the experimental data are in very reasonable agreement with this dependence (solid line); with increasing signal intensity by a factor of 4 the vortex radius decreases by a factor of 2. At small signal intensities the vortex radius does not diverge as expected from the $n_s^{-1/2}$, but is probably limited by the finite signal size.

Experimental constraints hinder the direct demonstration of the creation of the antivortex in the idler accompanying the signal vortex; the OPO idler is typically >100 times weaker than the signal [14] due to the small photonic component at the idler, making it impossible to obtain spatial and interference images. However, we are able to show that the pair scattering leads to the creation of vortex antivortex pairs. Using a pump power just below the OPO threshold, we apply a beam with OAM $M = +1$ at the angle and energy where the OPO idler would appear. The “signal” then forms at $k \sim 0$ due to pair scattering from the pump, stimulated by the probe beam. In this geometry, we are able to perform the required imaging of the signal; the vortex image and the corresponding self-interference image are shown in Figs. 1(f) and 1(g), respectively. A dip in intensity is observed in Fig. 1(f) corresponding to a

vortex core. In Fig. 1(g) a fork dislocation corresponding to a vortex labeled A in Fig. 1(f) is observed, pointing to the right rather than to the left as for the vortex labeled A in Figs. 1(d) and 1(e), proving that the OAM is $m_s = -1$. Here, in Fig. 1(f) we image the same original (uninverted) beam as in Fig. 1(d).

In order to obtain a better understanding of the nature of the vortex, we have developed a theoretical model of a vortex in a $\chi^{(3)}$ OPO. Our treatment starts from the classical model of an OPO with uniform fields for pump, signal, and idler [18], and adds a vortex antivortex pair to the signal and idler, which locally modulates the fields in the region of the core. The modulation functions are determined by a set of coupled nonlinear differential equations, which we solve numerically. Details of the theory will be presented elsewhere.

Figure 3(a) shows typical profiles calculated for the three modes at low ($\sqrt{n_s n_i} = 0.01 n_p$) and high ($\sqrt{n_s n_i} = n_p$) signal intensities. Here n_p is the polariton density in the pump, which is nearly constant above threshold [18]. The calculation uses structural parameters appropriate to the experimental microcavity. The length scale in the plot is the healing length $\xi = (2M_c \kappa \sqrt{n_s n_i})^{-1/2}$. At small pump powers when $\sqrt{n_s n_i} = 0.01 n_p$ the signal and idler profiles are identical with radii of $\sim \xi$, and the pump is unaffected by the vortex. At high powers, the signal and idler profiles separate, with the idler having the smaller core, and a “bump” appears in the pump in the core region, where it is less depleted by the OPO process. The separation is shown in more detail in Fig. 3(b), where the vortex radii are plotted as a function of $\sqrt{n_s n_i} / n_p$. In this figure, the length scaling used is $(2M_c \kappa n_p)^{-1/2}$, to show explicitly the strong dependence on population.

The behavior can be understood using physical arguments based on the competition between the blueshift due to polariton-polariton interactions and the OPO scattering process arising from the nonlinearity [18]. The blueshift provides restoring forces which control the vortex radius,

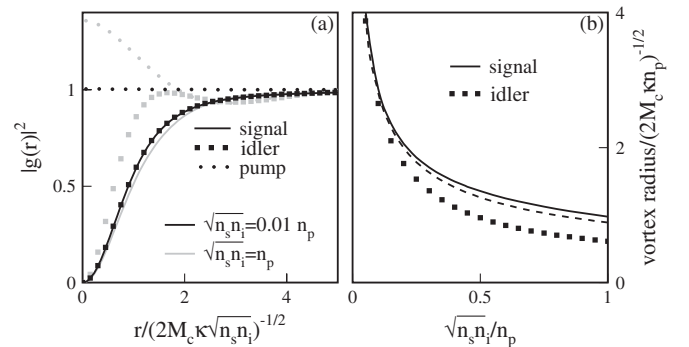


FIG. 3. (a) Calculated signal (solid line), idler (square), and pump (dot) vortex profiles for signal and idler populations near threshold (black), and comparable to the pump (gray). (b) Variation of signal (solid line) and idler (square) core radii with signal or idler population. The dashed line is obtained from the near-threshold approximation equation (1).

as in the standard Gross-Pitaevskii equation [8]. It tends, at high powers, to produce different radii for the signal and idler because of the different masses; the more excitonic idler has a larger mass and so a smaller core. The OPO effect, by contrast, tries to make the profiles the same, because the OPO scattering is a local process, generating equal numbers of signal and idler polaritons at a given point in space.

The regime when the signal (and idler) population is much smaller than that in the pump ($n_s \sim 10^9 - 10^{10} \text{ cm}^{-2} \leq 0.2n_p$), is the one which is accessed in the present experiment [14]. In this case the OPO process dominates and the signal and idler profiles are locked together. The actual profile, $g(r)$, is, however, still determined by the blueshift terms; we can show that it satisfies an equation of the complex Ginzburg-Landau (CGL) type. Scaling the radius by the healing length ξ , we find

$$\left(\nabla_r^2 - \frac{m^2}{r^2}\right)g = \alpha(|g|^2 - 1)g, \quad (1)$$

where α is a complex constant with magnitude of order unity. It is determined by the ratio of the signal and idler linewidths, the exciton and cavity amplitudes (Hopfield factors) for the modes, and the phase difference $\Phi = 2\phi_p - \phi_s - \phi_i$ between the pump, signal, and idler phases, respectively. The solutions of this equation provide the dashed curve on Fig. 3(b), which is in very good agreement with the full theory for $\sqrt{n_s n_i} \lesssim 0.1n_p$. We note that a CGL equation of the same form is obtained for the incoherently pumped polariton condensate [19]. Although the OPO is a more complicated system, with three coherent fields, in the near-threshold regime it displays similar physics.

In weak coupling microcavity lasers (VCSELs), which operate at 1 to 2 order of magnitude higher electron-hole density than the polariton condensate, there is also an effective $\chi^{(3)}$ nonlinearity due to saturation of the gain [7]. This contrasts with the direct interactions between polaritons due to their exciton component. At low powers the VCSEL can be approximated by a CGL type equation similar to Eq. (1), and hence a decrease of vortex size with power is expected. However, the variation is expected to be more complicated as saturation is approached. A typical vortex size of several microns is observed in Ref. [7], but the physical origin of this is not discussed. To the best of our knowledge there have not been any studies on the variation of vortex size with pump power in VCSELs. Instead as the current density was increased in Ref. [7], spontaneous vortex proliferation was observed.

It is also interesting to compare the magnitudes of the vortex radius in the polariton and atom systems: for polaritons, $M_c \sim 10^{-5}m_e$, and $\kappa\sqrt{n_s n_i} \sim 10^{-1} \text{ meV}$. Equivalent values for a rubidium BEC with a density $n \sim 10^{14} \text{ cm}^{-3}$ would be $M \sim 10^5 m_e$ and $gn = 4\pi a_s n \hbar^2 / M \sim 10^{-7} \text{ meV}$, where the atom scattering length $a_s \approx 5 \text{ nm}$ [8] and g is the atom interaction strength. This gives radii

for the polariton and the atom systems of order 10 and $0.1 \mu\text{m}$, respectively. We see that the much smaller polariton mass compensates the stronger interactions, giving a healing length 2 orders of magnitude greater than atom BECs.

To conclude, the polariton OPO supports a novel excitation consisting of a vortex in the signal and an antivortex in the idler. We have shown that the core radius of vortices in the polariton condensate is determined by polariton-polariton interactions. With ultrafast measurement techniques the imprinting method also provides a means to investigate vortex dynamics on time scales inaccessible in other systems. By including polarization as well as phase variation in the imprinting beam, it should be possible to study more complicated topological defects, such as ‘‘half-vortices’’ [20].

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